OC-seislet: Seislet transform construction with differential offset continuation

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ABSTRACT

Many of the geophysical data-analysis problems such as signal-noise separation and data regularization are conveniently formulated in a transform domain, in which the signal appears sparse. Classic transforms such as the Fourier transform or the digital wavelet transform (DWT) fail occasionally in processing complex seismic wavefields because of the nonstationarity of seismic data in time and space dimensions. We present a sparse multiscale transform domain specifically tailored to seismic reflection data. The new wavelet-like transform — the OC-seislet transform — uses a differential offset-continuation (OC) operator that predicts prestack reflection data in offset, midpoint, and time coordinates. It provides a high compression of reflection events. Its compression properties indicate the potential of OC seislets for applications such as seismic data regularization or noise attenuation. Results of applying the method to synthetic and field data examples demonstrate that the OC-seislet transform can reconstruct missing seismic data and eliminate random noise even in structurally complex areas.

INTRODUCTION

Digital wavelet transforms are excellent tools for multiscale data analysis. The wavelet transform is more powerful when compared to the classic Fourier transform because it is better fitted for representing nonstationary signals. Wavelets provide a sparse representation of piecewise regular signals, which can include transients and singularities (Mallat, 2009). In recent years, many wavelet-like transforms that explore directional characteristics of images (Starck et al., 2002; Do and Vetterli, 2005; Le Pennec and Mallat, 2005; Velisavljevic, 2005) have been proposed. The curvelet transform in particular has important applications in seismic imaging and data analysis (Douma and de Hoop, 2007; Chauris and Nguyen, 2008; Herrmann et al., 2008). Fomel (2006) and Fomel and Liu (2010) investigate the possibility of designing a wavelet-like transform tailored specifically to seismic data and introduce it as a “seislet transform.” Based on the digital wavelet transform (DWT), the seislet transform follows patterns of seismic events (such as local slopes in two dimensions and frequencies in one dimension) when analyzing those events at different scales. The seislet transform’s compression ability finds applications in common data-processing tasks such as data regularization and noise attenuation. However, the problem of pattern detection limits its use in further applications. In two dimensions, conflicting slopes at a single data point are difficult to detect reliably even when using advanced methods (Fomel, 2002). It is also difficult to estimate local slopes in the presence of strong noise. A similar situation occurs in the 1D case, in which it is difficult to exactly represent a known seismic signal using a limited set of frequencies.

Offset continuation (OC) is a process of seismic data transformation between different offsets (Deregowski and Rocca, 1981; Bolondi et al., 1982; Salvador and Savelli, 1982; Fomel, 2003c). Different types of dip moveout (DMO) operators (Hale, 1991) can be seen as a continuation to zero offset and derived as solutions to initial-value problems with the OC differential equation. In the shot-record domain, OC transforms to be shot continuation (Bagaini and Spagnolini, 1996; Spagnolini and Opreni, 1996; Fomel, 2003b). The 3D analog is known as azimuth moveout (AMO); (Biondi et al., 1998; Fomel, 2003a). Bleistein and Jaramillo (2000) develop a general platform for Kirchhoff data mapping, which includes OC as a special case.

In this paper, we propose incorporating OC as the prediction operator into the seislet transform. We design the transform in the log-stretch frequency domain, in which each frequency slice can be processed independently and in parallel. We expect the new seislet transform to perform better than the previously proposed seislet transform by plane-wave destruction (PWD) — the PWD-seislet transform (Fomel and Liu, 2010) — in cases of moderate velocity.
variations and complex structures that generate conflicting dips in the data.

THEORETICAL BASIS

OC-seislet transform follows the same construction strategy — the lifting scheme — as the PWD-seislet transform but uses the OC operator for prediction at each transform scale.

The lifting scheme for the seislet transform

The lifting scheme (Sweldens, 1995) provides a convenient approach for designing DWTs. The general procedure is as follows:

1) Organize the input data as a sequence of records. For OC-seislet transform of 2D seismic reflection data, the input is in the frequency-midpoint wavenumber-offset domain after the log-stretched NMO correction (Bolondi et al., 1982) and the transform direction is offset.

2) Divide the data records (along the offset axis in the case of the OC-seislet transform) into even and odd components \( e \) and \( o \), respectively. This step works at one scale level.

3) Find the residual difference \( r \) between the odd component and its prediction from the even component

\[
\mathbf{r} = \mathbf{o} - \mathbf{P}[\mathbf{e}],
\]

where \( \mathbf{P} \) is a prediction operator. For example, one can obtain Cohen-Daubechies-Feauveau (CDF) 5/3 biorthogonal wavelets (Cohen et al., 1992) by defining the prediction operator as a linear interpolation between two neighboring samples

\[
\mathbf{P}[\mathbf{e}]_k = (\mathbf{e}_{k-1} + \mathbf{e}_k)/2,
\]

where \( k \) is an index number at the current scale level.

4) Find an approximation \( \mathbf{c} \) of the data by updating the even component

\[
\mathbf{c} = \mathbf{e} + \mathbf{U}[\mathbf{r}],
\]

where \( \mathbf{U} \) is an update operator. Constructing the update operator for CDF 5/3 biorthogonal wavelets aims at preserving the running average of the signal (Sweldens and Schröder, 1996)

\[
\mathbf{U}[\mathbf{r}]_k = (\mathbf{r}_{k-1} + \mathbf{r}_k)/4.
\]

5) The coarse approximation \( \mathbf{c} \) becomes the new data and the sequence of steps is repeated on the new data to calculate the transform coefficients at a coarser scale level.

Next, we define new prediction and update operators using OC operators.

OC-seislet structure

We define the OC-seislet transform by specifying prediction and update operators with the help of the OC operator. Prediction and update operators for the OC-seislet transform are specified by modifying the biorthogonal wavelet construction in equations 2 and 4 as follows (Fomel, 2006; Fomel and Liu, 2010):

\[
\mathbf{P}[\mathbf{e}]_k = (\mathbf{S}_k^+[\mathbf{e}_{k-1}] + \mathbf{S}_k^-[\mathbf{e}_k])/2,
\]

\[
\mathbf{U}[\mathbf{r}]_k = (\mathbf{S}_k^+[\mathbf{r}_{k-1}] + \mathbf{S}_k^-[\mathbf{r}_k])/4,
\]

where \( \mathbf{S}_k^+[\cdot] \) and \( \mathbf{S}_k^-[\cdot] \) are operators that predict the data record (a common-offset section) by differential OC from its left- and right-neighboring common-offset sections with different offsets. Offset continuation operators provide the physical connection between data records. The theory of OC is reviewed in Appendix A.

One can also use a higher-order transform, for example, by using the template of the CDF 9/7 biorthogonal wavelet transform, which is used in JPEG2000 compression (Lian et al., 2001). There is only one stage (one prediction and one update) for the CDF 5/3 wavelet transform but there are two cascaded stages and one scaling operation for CDF 9/7 wavelet transform. Prediction and update operators for a high-order OC-seislet transform are defined as follows:

\[
\mathbf{P}_1[\mathbf{e}]_k = (\mathbf{S}_k^+[\mathbf{e}_{k-1}] + \mathbf{S}_k^-[\mathbf{e}_k])\cdot \alpha,
\]

\[
\mathbf{U}_1[\mathbf{r}]_k = (\mathbf{S}_k^+[\mathbf{r}_{k-1}] + \mathbf{S}_k^-[\mathbf{r}_k])\cdot \beta,
\]

\[
\mathbf{P}_2[\mathbf{e}]_k = (\mathbf{S}_k^+[\mathbf{e}_{k-1}] + \mathbf{S}_k^-[\mathbf{e}_k])\cdot \gamma,
\]

\[
\mathbf{U}_2[\mathbf{r}]_k = (\mathbf{S}_k^+[\mathbf{r}_{k-1}] + \mathbf{S}_k^-[\mathbf{r}_k])\cdot \delta,
\]

where subscripts 1 and 2 represent the first and second stages, respectively. Here, \( \alpha, \beta, \gamma, \) and \( \delta \) are defined numerically as follows:

\[
\alpha = -1.586134342,
\]

\[
\beta = -0.052980118,
\]

\[
\gamma = 0.882911076,
\]

\[
\delta = 0.443506852.
\]

One can combine equations 1, 3, 7, and 8 to finish the first stage and repeatedly process the result by using equations 1, 3, 9, and 10. The scale normalization factors correspond to the CDF 9/7 biorthogonal wavelet transform (Daubechies and Sweldens, 1998). Scaling and coefficients are as follows:

\[
\mathbf{e} = \mathbf{e} \cdot K,
\]

\[
\mathbf{o} = \mathbf{o} \cdot (1/K),
\]

where \( K = 1.230174105 \).

We used the high-order version of OC-seislet transform to process the synthetic and field data examples used in this paper.

Simple example

Figure 1 shows a 2D slice out of the benchmark French model (French, 1974). We created a 2D prestack data set (Figure 2a) by Kirchhoff modeling. Three sections in Figure 2a show the time slice at time position 0.6 s (top section), common-offset section at offset position of 0.2 km (bottom left section), and common-midpoint gather at midpoint position of 1.0 km (bottom right section). The reflector with a round dome and corners creates complicated reflection events along midpoint and offset axes. The inflection points of the reflector lead to traveltime triplications at some offsets. Figure 2b shows a preprocessed data cube in the \( F-K \) offset domain after the
log-stretched NMO correction and a double Fourier transform along the stretched time and midpoint axes. We apply the OC-seislet transform described above along the offset axis in Figure 2b. Thus, the offset axis becomes the scale axis. The cube of the transform coefficients is shown in Figure 3b and should be compared with the corresponding Fourier transform along the offset direction in Figure 3a. The OC-seislet transform coefficients get concentrated at small scales, which enable an effective compression. In contrast, the Fourier transform develops large coefficients at coarser scales but has small residual coefficients at fine scales. Figure 4 shows a comparison between the decay of coefficients (sorted from large to small) between the Fourier transform and the OC-seislet transform. A significantly faster decay of the OC-seislet coefficients is evident.

Iterative soft thresholding

The proposed OC-seislet transform uses physical OC to compress the reflection data after NMO and log-stretch transform of the time coordinate, followed by double Fourier transforms of the stretched time axis and midpoint axis. If seismic traces in the midpoint direction are missing, the Fourier transform can produce artifacts (spatial leakage) along the midpoint-wavenumber axis (Zwartjes and Gisolf, 2007). Additionally, the missing seismic traces in the offset direction affect the continuity of predicted data. Once the OC-seislet transform is applied for predicting and compressing, the artifacts, discontinuity information, and random noise spread over different scales while the predictable reflection information gets compressed to large coefficients at small scales. A simple thresholding method can easily remove the small coefficients of artifacts. Finally, applying the inverse OC-seislet transform, inverse FFTs in time and midpoint axes, inverse log stretch, and inverse NMO reconstructs the data while attenuating random noise, reducing artifacts along the midpoint axis, reconstructing continuity in the offset axis, and recovering main structural features without using any assumptions about structural continuity in the midpoint-offset-time domain. The key steps are shown in Figure 5a.

Figure 1. Two-dimensional slice out of the benchmark French model (French, 1974).

Figure 2. Two-dimensional synthetic prestack data in (a) the t-x-offset domain and in (b) the F-K-offset domain.

The idea of sparse transforms has been thoroughly explored in the literature, with application to Fourier and curvelet transforms. Liu and Sacchi (2004) propose a Fourier-based minimum-weighted norm interpolation (MWNI) algorithm with iterative inversion to perform multidimensional reconstruction of seismic wavefields. Xu et al. (2005) suggest an iterative Fourier-based matching pursuit (antileakage Fourier transform) for seismic data regularization. Abma and Kabir (2006) use a Fourier-based method with iterative thresholding to solve seismic data-interpolation problems. Herrmann and Hennenfent (2008) present a recovery method that exploits the curvelet frame. Missing data interpolation is a particular case of data regularization, in which the input data are already given on a regular grid and one needs to reconstruct only the missing values in empty bins (Fomel, 2001). For input data with nonuniform spatial sampling, one can bin the data to a regular grid first and then use the proposed method for filling empty bins. It is also possible to generalize the method for combining irregular data interpolation with a sparse...
transform (Zwartjes and Gisolf, 2007). The thresholding iteration helps a sparse transform recover the missing information. We adopt a similar thresholding strategy (fixed thresholding value for each iteration) with the OC-seislet transform. The key steps of the algorithm are illustrated schematically in Figure 5b.

SYNTHETIC DATA TESTS

We continue to use the synthetic model from Figure 1 to test the proposed method. Figure 6 illustrates the data with normally distributed random noise added. Figure 7a shows the data cube in the f-k-offset domain after the log-stretched NMO correction and a double Fourier transform along the stretched time axis and midpoint axis. We run the OC-seislet transform in parallel on individual frequency slices. Figure 7b shows the OC-seislet coefficients. All reflection information is concentrated in a small-scale range. However, because random noise cannot be predicted well by the OC opera-

Figure 3. (a) Fourier transform and (b) OC-seislet transform of the input data from Figure 2b along the offset axis.

Figure 4. Transform coefficients sorted from large to small, normalized, and plotted on a decibel scale.

Figure 5. Schematic illustration of the thresholding workflow. (a) Denoising with simple soft thresholding and (b) missing data interpolation with iterative soft thresholding.

Figure 6. Two-dimensional noisy data in the t-x-offset domain.
tor, it spreads throughout the whole transform domain. The inverse Fourier transform in time and midpoint directions and the inverse log stretch return the OC-seislet coefficients to the time-midpoint-scale domain. The coefficients at the zero scale represent stacking along the offset direction, which is equivalent to the DMO stack (Figure 8b).

We use a soft-thresholding method to separate reflection and random noise in the OC-seislet domain. Figure 8a displays the result after the inverse OC-seislet transform. Compared to Figure 7a, data in the $f$-$k$-offset domain contain only useful information. Figure 9b shows the denoising result in the $t$-$x$-offset domain after the double inverse Fourier transform, the inverse log stretch, and the inverse NMO. All characteristics of reflection and diffraction events are preserved well. For comparison, we used the PWD-seislet transform and soft-thresholding method with the same threshold values to process the noisy data (Figure 6). The result is shown in Figure 9a.

![Figure 7. (a) Noisy data in the $f$-$k$-offset domain and (b) OC-seislet coefficients in the $f$-$k$-scale domain.](image1)

![Figure 8. (a) Thresholded data in the $f$-$k$-offset domain and (b) OC-seislet coefficients in the $t$-$x$-scale domain.](image2)
Because the PWD-seislet transform is based on the local dip information, a mixture of different dips from the triplications makes it difficult to process the data in individual common-midpoint gathers. The PWD-seislet transform compresses all information along the local-event slopes in each common-midpoint gather. It separates the reflection signal and noise but smears the crossing events, especially at the far offset. The corresponding signal-to-noise ratios for denoised results with PWD-seislet transform and OC-seislet transform are 24.95 dB and 41.45 dB, respectively. The differences (Figure 10) between noisy data (Figure 6) and denoised results with the PWD-seislet transform (Figure 9a) and the OC-seislet transform (Figure 9b) further illustrate the effectiveness of the OC-seislet transform.

For a data-regularization test, we remove 80% of randomly selected traces (Figure 11a) from the ideal data (Figure 2a). The complex dip information makes it extremely difficult to interpolate the data in

Figure 9. Denoised result by different methods. (a) PWD-seislet transform and (b) OC-seislet transform (cf., Figure 2.).

Figure 10. Difference between noisy data (Figure 6) and denoised results with different methods (Figure 9). (a) PWD-seislet transform and (b) OC-seislet transform.
individual common-offset gathers. The data set is also nonstationary in the offset direction; therefore, a simple offset interpolation scheme would also fail. Figure 11b shows the data after the NMO correction, log-stretch transform, and double Fourier transforms. The missing traces introduce spatial artifacts in midpoint-wavenumber axis and discontinuities along the offset direction. After the OC-seislet transform, the reflection information can be predicted and compressed. Meanwhile, the artifacts spread over the whole transform domain (Figure 12a). The simple soft-thresholding algorithm (i.e., the iterative strategy with only one iteration) removes most of the artifacts and the inverse OC-seislet transform reconstructs the major reflection information according to the OC prediction (Figure 12b).

One can also use an iterative soft-thresholding strategy to implement missing data interpolation. This method recovers missing traces as long as seismic data are sparse enough in the transform domain.
To demonstrate the superior sparseness of the OC-seislet coefficients, we compare the proposed method with the 3D Fourier transform. Figure 13a displays the interpolated result after a 3D Fourier interpolation using iterative thresholding (Abma and Kabir, 2006). The Fourier transform cannot provide enough sparseness of coefficients for complex reflections and therefore fails in recovering all missing traces. The OC-seislet transform is based on a physical prediction and provides a much sparser domain for reflections and diffractions. Iterative thresholding succeeds in interpolating the missing traces (Figure 13b).

FIELD DATA TESTS

We used a historic marine data set from the Gulf of Mexico (Claerbout, 2000) to evaluate the proposed method for noise attenuation and missing data-interpolation problems. The input data are shown in Figure 14a. Near- and far-offset information is completely missing. After log-stretched NMO and Fourier transforms along time and midpoint axes convert the original data from time-midpoint-offset domain to frequency-wavenumber-offset domain, the OC-seislet transform compresses the predictable information according to the

![Figure 13](image13.png)

Figure 13. Interpolated results using iterative thresholding with different sparse transforms. (a) Three-dimensional Fourier transform and (b) OC-seislet transform (cf., Figure 2).

![Figure 14](image14.png)

Figure 14. (a) Marine data in the t-x-offset domain and (b) OC-seislet coefficients in the f-k-scale domain.
physical connection between offsets. Unpredictable artifacts and reflections display different tendencies: Artifacts disperse throughout the whole transform domain while reflections are compressed to a small-scale range (Figure 14b). If we choose the significant coefficients using soft thresholding, the inverse processing steps (including the inverse OC-seislet transform, the inverse FFTs in time and mid-point axes, and the inverse log-stretch NMO) effectively remove incoherent artifacts from the data (Figure 15a). The difference between the input data (Figure 14a) and the denoised result (Figure 15a) is shown in Figure 15b. Most coherent events including reflections and diffractions are well protected by the proposed method thanks to the ability of the OC operator to predict structurally complex wavefields. However, some energy of the events is partly missing. To achieve a better result, one might need to use a more accurate NMO velocity. Thus, denoising is a naturally defined operation in the OC-seislet domain. The soft-thresholding techniques eliminate only part of spatial aliasing because of sparse offset sampling. Additional offset interpolation can further remove spatial aliasing.

For a data regularization test, we removed 70% of randomly selected traces from the input data (Figure 16). After applying the OC-seislet transform and iterative soft thresholding, spatial artifacts and offset discontinuity in the transform domain coming from missing traces get attenuated. The interpolated result is shown in Figure 17b. Missing traces have been interpolated well even when diffractions are present. For comparison, we applied 3D Fourier transform with iterative soft thresholding as proposed by Abma and Kabir (2006).

The same parameters of soft thresholding for those in the OC-seislet transform were used. The final result is shown in Figure 17a. The Fourier transform fails to recover the missing information as accurately as the OC-seislet transform. Fundamentally, the Fourier coefficients are not sparse enough for this data set. Figure 18 shows the comparison after DMO stacking. The result from 3D Fourier transform displays obvious gaps that are caused by imperfect interpolation (Figure 18a) whereas the OC-seislet transform and iterative

Figure 15. (a) Denoised result in the $t$-$x$-offset domain and (b) difference between the input data (Figure 14a) and the denoised result (Figure 15a).

Figure 16. Marine data with 70% traces removed.

Figure 17. Interpolated results using iterative thresholding with different transforms. (a) Three-dimensional Fourier transform and (b) OC-seislet transform.
We thank Tamas Nemeth, Mauricio Sacchi, Sandra Tegtmeier-Last, and two anonymous reviewers for their constructive comments and suggestions. This publication was authorized by the Director, Bureau of Economic Geology, The University of Texas at Austin.

APPENDIX A

REVIEW OF DIFFERENTIAL OFFSET CONTINUATION

In this appendix, we review the theory of differential OC from Fomel (2003a, c). The partial differential equation for OC (differential AMO) takes the form

\[
\left(\frac{\partial^2 P}{\partial h^2} - h^2 \frac{\partial^2 P}{\partial t^2}\right) = h \frac{\partial^2 P}{\partial t \partial h}.
\]

(A-1)

where \( P(x, h, t) \) is the seismic data in the midpoint-offset-time domain, \( t \) is the time coordinate after the NMO correction, \( h^T \) denotes the transpose of \( h \), and \( P_{zz} \) is the tensor of the second-order midpoint derivatives.

A particularly efficient implementation of OC results from a log-stretch transform of the time coordinate (Bolondi et al., 1982), followed by a Fourier transform of the stretched time axis. After these transforms, the OC equation takes the form

\[
\left(\frac{\partial^2 \tilde{P}}{\partial h^2} - \frac{\partial^2 \tilde{P}}{\partial \xi^2}\right) = i\Omega \frac{\partial \tilde{P}}{\partial \xi}.
\]

(A-2)

where \( \Omega \) is the dimensionless frequency corresponding to the stretched time coordinate and \( \tilde{P}(x, h, \Omega) \) is the transformed data. As in other frequency-space methods, equation A-2 can be applied independently and in parallel on different frequency slices.

In the frequency-wavenumber domain, the extrapolation operator is defined by solving an initial-value problem for equation A-2. The analytical solution takes the form

\[
\tilde{P}(k, h, \Omega) = \tilde{P}(k, h, \Omega) \frac{Z_d(k \cdot h_2)}{Z_d(k \cdot h_1)},
\]

(A-3)

where \( \tilde{P}(k, h, \Omega) \) is the double Fourier-transformed data, \( \lambda = (1 + i\Omega)/2 \), \( Z_d \) is the special function defined as

\[
Z_d(x) = I(1 - \lambda) \left(\frac{x}{2}\right)^\lambda J_{-\lambda}(x) =_0F_1\left(1 - \lambda; -\frac{x^2}{4}\right),
\]

(A-4)

where \( I \) is the gamma function, \( J_{-\lambda} \) is the Bessel function, and \( _0F_1 \) is the confluent hypergeometric limit function (Petkovsek et al., 1996). The wavenumber \( k \) in equation A-3 corresponds to the midpoint \( x \) in the original data domain. In high-frequency asymptotics, the OC operator takes the form

\[
\tilde{P}(k, h, \Omega) = \tilde{P}(k, h, \Omega) \frac{\exp\left[i\Omega \psi\left(\frac{2k \cdot h_2}{\Omega}\right)\right]}{\exp\left[i\Omega \psi\left(\frac{2k \cdot h_1}{\Omega}\right)\right]},
\]

(A-5)

where

\[
\psi(z) = \frac{z}{2} - \frac{\ln(z)}{2}.
\]
\[ F(\epsilon) = \sqrt{\frac{1 + \sqrt{1 + \epsilon^2}}{2 \sqrt{1 + \epsilon^2}}} \exp\left(\frac{1 - \sqrt{1 + \epsilon^2}}{2}\right). \]  
(A-6)

and

\[ \psi(\epsilon) = \frac{1}{2} \left(1 - \sqrt{1 + \epsilon^2} + \ln\left(1 + \sqrt{1 + \epsilon^2}\right)\right). \]  
(A-7)

The phase function \( \psi \) defined in equation A-7 corresponds to the analogous term in the exact-log DMO and AMO (Liner, 1990; Zhou et al., 1996; Biondi and Vlad, 2002).

REFERENCES


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